

Foundations of Quantum Mechanics (APM 421 H)

Winter 2014 Solutions 3 (2014.09.26)

Hilbert Spaces & Operators

Homework Problems

8. Orthogonal subspaces and projections onto subspaces (16 points)

Let $\{\varphi_n\}_{n\in\mathbb{N}}$ be an orthonormal basis (ONB) of a Hilbert space \mathcal{H} and $N\in\mathbb{N}$.

- (i) Prove that $E := \{\varphi_1, \dots, \varphi_N\}^{\perp}$ is a sub vector space.
- (ii) Give an ONB for the subspace $E = \{\varphi_1, \dots, \varphi_N\}^{\perp}$.
- (iii) Show that $(\{\varphi_1, \ldots, \varphi_N\}^{\perp})^{\perp} = E^{\perp} = \operatorname{span}\{\varphi_1, \ldots, \varphi_N\}.$

Moreover, define the map

$$P: \mathcal{H} \longrightarrow \mathcal{H}, \ P\psi := \sum_{n=1}^{N} \langle \varphi_n, \psi \rangle \ \varphi_n$$

- (iv) Show that P is linear, i. e. for any $\varphi, \psi \in \mathcal{H}$ and $\alpha \in \mathbb{C}$, we have $P(\alpha \varphi + \psi) = \alpha P \varphi + P \psi$.
- (v) Show that *P* is a projection, i. e. $P^2 = P$.
- (vi) Show that P is bounded, i. e. $||P\varphi|| \le ||\varphi||$ holds for any $\varphi \in \mathcal{H}$.

Solution:

(i) The orthogonal complement is defined as

$$E \stackrel{[1]}{=} \left\{ \psi \in \mathcal{H} \mid \langle \varphi_j, \psi \rangle = 0, \ j = 1, \dots, N \right\}.$$

For any $\phi, \psi \in E$ and $\alpha \in \mathbb{C}$, also the vector $\alpha \phi + \psi$ is an element of E [1]: for all j = 1, ..., N

$$\langle \varphi_j, \alpha \phi + \psi \rangle = \alpha \ \langle \varphi_j, \phi \rangle + \langle \varphi_j, \psi \rangle \stackrel{[1]}{=} 0$$

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is satisfied. Hence, E is a linear subspace of \mathcal{H} .

(ii) $\{\varphi_j\}_{j=N+1}^{\infty}$ [1]

(iii) Then $\varphi_i \in E^{\perp}$, because by definition of E

 $\langle \varphi_j, \psi \rangle \stackrel{[1]}{=} 0$

holds for all $\psi \in \mathcal{H}$. Thus, $\varphi_j \in E^{\perp}$ for all $j = 1, \ldots, N$. By (i), E is a linear sub space [1].

Now assume that there exists a $\psi \in E^{\perp}$ which is *not* a linear combination of $\{\varphi_1, \ldots, \varphi_N\}$ [1]. Since $\{\varphi_j\}_{j\in\mathbb{N}}$ is an orthonormal basis of \mathcal{H} , we can express ψ as

$$\psi = \sum_{j=1}^{\infty} c_j \,\varphi_j \,. \tag{1}$$

By assumption, there exists a $n \ge N + 1$ for which $c_n \ne 0$ [1]. But then

$$\langle \varphi_n, \psi \rangle = c_n \neq 0$$

and ψ cannot be an element of $E^{\perp},$ contradiction [1].

Hence, $E^{\perp} = \operatorname{span}\{\varphi_1, \ldots, \varphi_N\}.$

(iv) The linearity of P follows from the linearity of the scalar product in the first argument: for all $\phi, \psi \in \mathcal{H}$ and $\alpha \in \mathbb{C}$, we have

$$P(\alpha \phi + \psi) \stackrel{[1]}{=} \sum_{j=1}^{N} \langle \varphi_j, \alpha \phi + \psi \rangle \varphi_j = \alpha \sum_{j=1}^{N} \langle \varphi_j, \phi \rangle \varphi_j + \sum_{j=1}^{N} \langle \varphi_j, \psi \rangle \varphi_j$$
$$\stackrel{[1]}{=} \alpha P \phi + P \psi.$$

Hence, P is linear.

(v) For any $\psi \in \mathcal{H}$, we deduce using the linearity of *P*:

$$P^{2}\psi = P\left(\sum_{j=1}^{N} \langle \varphi_{j}, \psi \rangle \ \varphi_{j}\right) \stackrel{[1]}{=} \sum_{k=1}^{N} \langle \varphi_{j}, \psi \rangle \ P\varphi_{j}$$
$$= \sum_{k,j=1}^{N} \langle \varphi_{j}, \psi \rangle \underbrace{\langle \varphi_{k}, \varphi_{j} \rangle}_{=\delta_{k,j}} \varphi_{k} = \sum_{j=1}^{N} \langle \varphi_{j}, \psi \rangle \ \varphi_{j} \stackrel{[1]}{=} P\psi$$

Hence, P is a projection.

(vi) With the help of Bessel's inequality [1], we obtain the claim:

$$\left\|P\psi\right\| = \left\|\sum_{j=1}^{N} \left\langle\varphi_{j},\psi\right\rangle \varphi_{j}\right\| \stackrel{[1]}{\leq} \left\|\psi\right\|$$

9. The Fock space (13 points)

Let \mathcal{H} be a separable Hilbert space and $\mathcal{H}^{\otimes n} = \mathcal{H} \otimes \cdots \otimes \mathcal{H}$ the *n*-fold tensor product. By definition we set $\mathcal{H}^{\otimes 0} := \mathbb{C}$. Define the *Fock space* over \mathcal{H} as $\mathfrak{F}(\mathcal{H}) := \bigoplus_{n=0}^{\infty} \mathcal{H}^{\otimes n}$ with scalar product

$$\langle \varphi, \psi \rangle_{\mathfrak{F}} := \sum_{n=0}^{\infty} \langle \varphi_n, \psi_n \rangle_{\mathcal{H}^{\otimes n}}, \qquad \varphi = (\varphi_0, \varphi_1, \ldots), \psi = (\psi_0, \psi_1, \ldots) \in \mathfrak{F}(\mathcal{H}).$$

In the physics literature, $\mathcal{H}^{\otimes n}$ is called the *n*-particle Fock sector.

- (i) Give a countable orthonormal basis of $\mathfrak{F}(\mathcal{H})$. (No proof is necessary.)
- (ii) Show that for any element $\psi = (\psi_0, \psi_1, \psi_2, ...)$ of $\mathfrak{F}(\mathcal{H})$

$$\lim_{n \to \infty} \left\| \psi_n \right\|_{\mathcal{H}^{\otimes n}} = 0$$

holds where ψ_n is the corresponding element of the *n*-particle Fock sector.

Let A be a linear, bounded operator on \mathcal{H} with domain $\mathcal{D}(A) = \mathcal{H}$. Then we define the *second* quantization of A as the operator

$$\mathsf{d}\Gamma(A) := \bigoplus_{n=0}^{\infty} (A \otimes \mathsf{id}_{\mathcal{H}} \otimes \cdots \otimes \mathsf{id}_{\mathcal{H}} + \ldots + \mathsf{id}_{\mathcal{H}} \otimes \cdots \otimes \mathsf{id}_{\mathcal{H}} \otimes A)$$

acting on $\mathfrak{F}(\mathcal{H})$.

- (iii) Give the action of $d\Gamma(A)$ on the *n*th Fock sector $\mathcal{H}^{\otimes n}$, i. e. find $(d\Gamma(A)\psi)_n$.
- (iv) Give the domain of $d\Gamma(id_{\mathcal{H}})$ and discuss the physical meaning of $d\Gamma(id_{\mathcal{H}})$.

Solution:

(i) Since \mathcal{H} is countable, there exists a countable orthonormal basis $\{\varphi_n\}_{n\in\mathbb{N}}$ [1]. Then

$$\left\{\left(1,\varphi_{k_{1\,1}},\varphi_{k_{2\,1}}\otimes\varphi_{k_{2\,2}},\ldots,\varphi_{k_{n\,1}}\otimes\cdots\otimes\varphi_{k_{n\,n}},\ldots\right)\right\}_{\substack{n\in\mathbb{N}\\1\leq j\leq n\\k_{n\,j}\in\mathbb{N}}}$$
[2]

is a countable orthonormal basis of $\mathfrak{F}(\mathcal{H})$.

(ii) Pick an arbitrary $\psi = (\psi_0, \psi_1, \ldots) \in \mathfrak{F}(\mathcal{H})$. Then

$$\|\psi\|^2 = \sum_{n=0}^{\infty} \|\psi_n\|^2_{\mathcal{H}^{\otimes n}} < \infty$$
^[1]

which necessarily means $\|\psi_n\|_{\mathcal{H}^{\otimes n}} \to 0$ as $n \to \infty$ [1].

(iii)
$$(\mathrm{d}\Gamma(A)\psi)_n = (A \otimes \mathrm{id}_{\mathcal{H}} \otimes \cdots \otimes \mathrm{id}_{\mathcal{H}} + \ldots + \mathrm{id}_{\mathcal{H}} \otimes \cdots \otimes \mathrm{id}_{\mathcal{H}} \otimes A)\psi_n$$
 [1]

(iv) The second quantization of the identity is the so-called (particle) number operator [1] which – as the name suggests – measures the number of particles. The operator acts on the *n*th Fock sector $\mathcal{H}^{\otimes n}$ as

$$\underbrace{\mathsf{id}_{\mathcal{H}}\otimes\mathsf{id}_{\mathcal{H}}\otimes\cdots\otimes\mathsf{id}_{\mathcal{H}}+\ldots+\mathsf{id}_{\mathcal{H}}\otimes\cdots\otimes\mathsf{id}_{\mathcal{H}}\otimes\mathsf{id}_{\mathcal{H}}}_{n\,\mathsf{terms}} \stackrel{[\underline{1}]}{=} n\,\mathsf{id}_{\mathcal{H}^{\otimes n}}$$

and hence, we have

$$\mathrm{d}\Gamma(\mathrm{id}_{\mathcal{H}})\stackrel{[2]}{=}\bigoplus_{n=0}^{\infty}n\,\mathrm{id}_{\mathcal{H}^{\otimes n}}.$$

The domain consists of states with a finite number of particles [1], i. e.

$$\begin{split} \mathcal{D}\big(\mathsf{d}\Gamma(\mathsf{id}_{\mathcal{H}})\big) \stackrel{[1]}{=} \Big\{\varphi \in \mathfrak{F}(\mathcal{H}) \ \Big| \ \mathsf{d}\Gamma(\mathsf{id}_{\mathcal{H}})\varphi \in \mathfrak{F}(\mathcal{H}) \Big\} \\ \stackrel{[1]}{=} \Big\{\varphi = (\varphi_0, \varphi_1, \ldots) \in \mathfrak{F}(\mathcal{H}) \ \Big| \ \sum_{n=0}^{\infty} n^2 \big| \big\langle \varphi_n, \varphi_n \big\rangle_{\mathcal{H}^{\otimes n}} \big|^2 < \infty \Big\}. \end{split}$$

10. Best approximation: Fourier series (12 points)

Let $L^2([0, 2\pi])$ be the Hilbert space of square-integrable functions with scalar product

$$\langle f,g \rangle := \frac{1}{2\pi} \int_0^{2\pi} \mathrm{d}x \,\overline{f(x)} \,g(x)$$

- (i) Show that $\{e^{+inx}\}_{n\in\mathbb{Z}}$ is an orthonormal system of vectors in $L^2([0, 2\pi])$.
- (ii) Let $E := \{e^{+inx}\}_{\substack{n \in \mathbb{Z} \\ |n| \leq 4}}$, and consider the functions $f(x) = \sin(2x)$ and g(x) = x. Give the element of best approximation of f and g in span E.
- (iii) Why don't these arguments work for $L^2(\mathbb{R})$?

Solution:

(i) The scalar product for $n \neq m$ is

$$\begin{split} \left\langle \mathbf{e}^{+\mathrm{i}nx}, \mathbf{e}^{+\mathrm{i}mx} \right\rangle &\stackrel{[1]}{=} \frac{1}{2\pi} \int_{0}^{2\pi} \mathrm{d}x \, \mathbf{e}^{-\mathrm{i}nx} \, \mathbf{e}^{+\mathrm{i}mx} \\ &= \frac{1}{2\pi} \int_{0}^{2\pi} \mathbf{e}^{+\mathrm{i}(m-n)x} = \frac{1}{2\pi} \frac{1}{i(m-n)} \, \mathbf{e}^{+\mathrm{i}(m-n)x} \Big|_{0}^{2\pi} \stackrel{[1]}{=} 0. \end{split}$$

For n = m we obtain instead

$$\langle \mathbf{e}^{+\mathrm{i}nx}, \mathbf{e}^{+\mathrm{i}mx} \rangle \stackrel{[1]}{=} \frac{1}{2\pi} \int_0^{2\pi} \mathrm{d}x \, \mathbf{e}^{+\mathrm{i}(m-n)x} = \frac{1}{2\pi} \int_0^{2\pi} \mathrm{d}x \, 1 \stackrel{[1]}{=} 1.$$

The vectors are all mutually orthogonal and are normed to 1, and hence, $\{e^{+inx}\}_{n\in\mathbb{Z}}$ is an orthonormal system [1].

(ii) Since $f(x) \stackrel{[1]}{=} \frac{1}{i2} (e^{+i2x} - e^{-2x})$ it lies in *E*, and *f* is its own Element of Best Approximation [1]. For *g* we proceed to compute the Fourier components, starting with n = 0:

$$c_0 = \frac{1}{2\pi} \int_0^{2\pi} \mathrm{d}x \, x \stackrel{[1]}{=} \pi$$

For $n \neq 0$ we obtain

$$c_n \stackrel{[1]}{=} \frac{1}{2\pi} \int_0^{2\pi} dx \, x \, e^{+inx} = \frac{1}{i2\pi n} x \, e^{+inx} \Big|_0^{2\pi} - \frac{1}{i2\pi n} \int_0^{2\pi} dx \, 1 \cdot e^{inx} \\ = \frac{1}{in} - \frac{1}{2\pi} \left[\frac{1}{(in)^2} \, e^{+inx} \right]_0^{2\pi} \stackrel{[1]}{=} \frac{1}{in}$$

Thus, the Element of Best Approximation of g in E is

$$g_E(x) = \pi + \sum_{n=1}^4 \frac{1}{in} \left(e^{+inx} - e^{-inx} \right) \stackrel{[1]}{=} \pi + \sum_{n=1}^4 \frac{2}{n} \sin nx.$$

(iii) The functions $e^{+inx} \notin L^2(\mathbb{R})$ are not square integrable [1], and hence, the arguments cannot be repeated.

11. Multiplication operators (14 points)

Let $V \in L^{\infty}(\mathbb{R}^d)$ and define the multiplication operator

$$(T_V\psi)(x) := V(x)\,\psi(x)\,,\qquad\qquad \psi \in L^2(\mathbb{R}^d)\,.$$

- (i) Show that $T_V: L^2(\mathbb{R}^d) \longrightarrow L^2(\mathbb{R}^d)$ is bounded.
- (ii) Assume $V \in L^{\infty}(\mathbb{R}^d)$ is real-valued. Show that then $\langle \varphi, T_V \psi \rangle_{L^2(\mathbb{R}^d)} = \langle T_V \varphi, \psi \rangle_{L^2(\mathbb{R}^d)}$ holds for all $\varphi, \psi \in L^2(\mathbb{R}^d)$, i. e. T_V is selfadjoint.
- (iii) Assume that V is bounded away from 0 and $+\infty$, i. e. that there exist C > c > 0 so that

$$0 < c \le V(x) \le C < +\infty$$

holds for all $x \in \mathbb{R}^d$. Find the inverse of T_V and show the inverse is bounded.

Solution:

(i) From the elementary estimate $|(T_V\psi)(x)| = |V(x)\psi(x)| \le ||V||_{\infty} |\psi(x)|$ [1], we deduce

$$\begin{split} \left\| T_V \psi \right\| &\stackrel{[1]}{=} \left(\int_{\mathbb{R}^d} \mathrm{d}x \left| (T_V \psi)(x) \right|^2 \right)^{1/2} \\ &\stackrel{[1]}{\leq} \left(\int_{\mathbb{R}^d} \mathrm{d}x \left\| V \right\|_{\infty}^2 |\psi(x)|^2 \right)^{1/2} \\ &= \| V \|_{\infty} \left(\int_{\mathbb{R}^d} \mathrm{d}x \left| \psi(x) \right|^2 \right)^{1/2} \\ &\stackrel{[1]}{=} \| V \|_{\infty} \| \psi \| \,. \end{split}$$

Hence, T_V is bounded [1].

(ii) The claim follows from $\overline{V} = V$ and direct computation: for any $\varphi, \psi \in L^2(\mathbb{R}^d)$, we have

$$\begin{split} \langle \varphi, T_V \psi \rangle \stackrel{[1]}{=} & \int_{\mathbb{R}^d} \mathrm{d}x \, \overline{\varphi(x)} \, (T_V \psi)(x) \stackrel{[1]}{=} \int_{\mathbb{R}^d} \mathrm{d}x \, \overline{\varphi(x)} \, V(x) \, \psi(x) \\ &= \int_{\mathbb{R}^d} \mathrm{d}x \, \overline{V(x) \, \varphi(x)} \, \psi(x) = \int_{\mathbb{R}^d} \mathrm{d}x \, \overline{(T_V \varphi)(x)} \, \psi(x) \\ & \stackrel{[1]}{=} \langle T_V \varphi, \psi \rangle \; . \end{split}$$

(iii) Since V is bounded away from $0 \text{ and } +\infty$, so is V^{-1} [1],

$$0 < C^{-1} \le V^{-1}(x) \le c^{-1} < \infty.$$

Hence, also $T_{V^{-1}}: L^2(\mathbb{R}^d) \longrightarrow L^2(\mathbb{R}^d)$ is a bounded multiplication operator by (i) [1]. Moreover, by direct computation, we verify that $T_{V^{-1}}$ is the inverse to T_V [1], e. g.

$$(T_V T_{V^{-1}} \psi)(x) \stackrel{[1]}{=} V(x) (T_{V^{-1}} \psi)(x)$$

= $V(x) V^{-1}(x) \psi(x) \stackrel{[1]}{=} \psi(x) ,$

and similarly $T_{V^{-1}} T_V = \operatorname{id}_{L^2(\mathbb{R}^d)} [1].$