

# Foundations of Quantum Mechanics (APM 421 H)

Winter 2014 Solutions 4 (2014.10.03)

# Operators

### **Homework Problems**

# 12. Projections (19 points)

Consider the multiplication operator  $P = p(\hat{x})$  on  $L^2(\mathbb{R}^d)$  associated to the function

$$p(x) = \begin{cases} 1 & x \ge 0\\ 0 & x < 0 \end{cases}$$

- (i) Find 2 eigenfunctions.
- (ii) Compute  $\sigma(P)$ .
- (iii) Determine the nature of the spectrum, i. e. determine  $\sigma_p(P)$ ,  $\sigma_{cont}(P)$ , and  $\sigma_r(P)$ .
- (iv) Prove that P is an orthogonal projection.

#### Solution:

(i) For instance, consider the functions

$$\psi_0(x) \stackrel{[1]}{=} \begin{cases} 1 & x \in [-1,0] \\ 0 & \text{else} \end{cases}, \qquad \qquad \psi_1(x) \stackrel{[1]}{=} \begin{cases} 1 & x \in [0,1] \\ 0 & \text{else} \end{cases}$$

(All that is important is that one of them is zero on  $[0, +\infty)$  and the other is zero on  $(-\infty, 0]$ .) Then clearly,  $\psi_0$  is an eigenfunction to the eigenvalue 0,  $(P\psi_0)(x) = p(x) \psi_0(x) = 0$  [1], while  $\psi_1$  is an eigenfunction to the eigenvalue 1,  $(P\psi_1)(x) = p(x) \psi_1(x) = \psi_1(x)$  [1].

(ii) First of all, P has two eigenvalues, namely 0 and 1 [2]: the eigenvectors to the eigenvalue 1 are functions which vanish almost everywhere on  $(-\infty, 0)$ . Similarly, eigenvectors to the eigenvalue 0 are functions which vanish almost everywhere on  $[0, +\infty)$ .

Since  $(P - z)\varphi = 0$  means that  $((P - z)\varphi)(x) = 0$  for almost all  $x \in \mathbb{R}$  [1]. Thus, for all  $z \neq 0, 1$  we have  $(P - z)\varphi \neq 0$  for all  $\varphi \neq 0$  [1], i. e. P - z is invertible as long as  $z \neq 0, 1$  [1], and we have shown  $\sigma(P) = \{0, 1\}$  [1].

(iii) Since all vectors are eigenvectors, the spectrum consists just of point spectrum,

$$\sigma_{\mathbf{p}}(P) \stackrel{[1]}{=} \sigma(P) = \{0, 1\}, \qquad \sigma_{\mathbf{cont}}(P) \stackrel{[1]}{=} \emptyset, \qquad \sigma_{\mathbf{r}}(P) \stackrel{[1]}{=} \emptyset.$$

(iv) p is a real-valued, bounded function, and hence, by problem 11 (ii)  $P^* = P$  [2]. (Otherwise, one needs to show this by hand for this special case.)

To see  $P^2 = P$ , we note  $p^2 = p$  in the sense of functions and conclude

$$(P^{2}\varphi)(x) \stackrel{[1]}{=} p(x)^{2} \varphi(x) \stackrel{[1]}{=} p(x) \varphi(x) \stackrel{[1]}{=} (P\varphi)(x).$$

Thus,  $P = P^* = P^2$  is an orthogonal projection [1].

#### 13. The discrete Laplacian (24 points)

Consider the Hilbert space  $\ell^2(\mathbb{Z})$  with the usual scalar product  $\langle \cdot, \cdot \rangle_{\ell^2(\mathbb{Z})}$ . Define the shift operator

$$\mathfrak{s}: \ell^2(\mathbb{Z}) \longrightarrow \ell^2(\mathbb{Z}), \ (\mathfrak{s}\psi)(n) := \psi(n-1)$$

as well as the shift by  $a \in \mathbb{Z}$  lattice units,  $\mathfrak{s}_a := \mathfrak{s}^a$ . Consider the discrete Laplacian

$$\Delta: \ell^2(\mathbb{Z}) \longrightarrow \ell^2(\mathbb{Z}), \ (\Delta\psi)(n):=\psi(n+1)+\psi(n-1)-2\psi(n)$$

- (i) Compute  $\mathfrak{s}_a^*$  and prove that  $\mathfrak{s}_a$  is unitary.
- (ii) Show that  $\Delta$  is a bounded operator on  $\ell^2(\mathbb{Z})$ .
- (iii) Show that  $\mathfrak{s}_a$  and  $\Delta$  commute, i. e.  $[\mathfrak{s}_a, \Delta] := \mathfrak{s}_a \Delta \Delta \mathfrak{s}_a = 0$ .
- (iv) Compute  $\Delta^*$ .
- (v) Determine  $E_k$  so that

$$\psi_k(n) := \mathbf{e}^{+\mathbf{i}nk}, \qquad n \in \mathbb{Z}, k \in [-\pi, +\pi],$$

is an eigenvalue to the discrete Laplacian,

$$(\Delta \psi_k)(n) = E_k \psi_k(n).$$

Is  $\psi_k$  an element of  $\ell^2(\mathbb{Z})$ ?

**Remark:** The Hilbert space  $\ell^2(\mathbb{Z})$  is often used in solid state physics where the shift operator  $(\widehat{\mathfrak{s}\psi})(n) := \widehat{\psi}(n-1)$  is interpreted as translating the particle by one lattice unit.

# Solution:

(i) Let  $\varphi, \psi \in \ell^2(\mathbb{Z})$  and  $a \in \mathbb{Z}$ . The adjoint operator  $\mathfrak{s}_a^*$  is then  $\mathfrak{s}_{-a}$ ,

$$\begin{split} \left\langle \varphi, \mathfrak{s}_{a} \psi \right\rangle \stackrel{[1]}{=} \sum_{n \in \mathbb{Z}} \overline{\varphi(n)} \, (\mathfrak{s}_{a} \psi)(n) \stackrel{[1]}{=} \sum_{n \in \mathbb{Z}} \overline{\varphi(n)} \, \psi(n-a) \stackrel{[1]}{=} \sum_{k \in \mathbb{Z}} \overline{\varphi(k+a)} \, \psi(k) \\ \stackrel{[1]}{=} \sum_{k \in \mathbb{Z}} \overline{(\mathfrak{s}_{-a} \varphi)(k)} \, \psi(k) \stackrel{[1]}{=} \left\langle \mathfrak{s}_{-a} \varphi, \psi \right\rangle. \end{split}$$

 $\mathfrak{s}_{-a}$  is also the inverse to  $\mathfrak{s}_a$  [1], since

$$(\mathfrak{s}_{-a}\mathfrak{s}_a\varphi)(n) = (\mathfrak{s}_a\varphi)(n+a) = \varphi(n+a-a) = \varphi(n)$$

holds for all  $\varphi \in \ell^2(\mathbb{Z})$  and  $n \in \mathbb{Z}$ . This means  $\mathfrak{s}_a$  is unitary [1].

(ii) We recognize that actually  $\Delta = \mathfrak{s} + \mathfrak{s}^* - 2$  [1], so we deduce  $\Delta$  is bounded

$$\|\Delta\| \stackrel{[1]}{\leq} \|\mathfrak{s}\| + \|\mathfrak{s}^*\| + \|2\| \stackrel{[1]}{=} 4.$$

(iii) Translations commute amongst one another,  $\mathfrak{s}_a \mathfrak{s}_b = \mathfrak{s}_{a+b} = \mathfrak{s}_b \mathfrak{s}_a$  [1], and hence,

$$\begin{split} \mathfrak{s}_a \Delta \stackrel{[1]}{=} \mathfrak{s}_a \, \mathfrak{s} + \mathfrak{s}_a \, \mathfrak{s}^* - 2 \mathfrak{s}_a \\ \stackrel{[1]}{=} \mathfrak{s} \, \mathfrak{s}_a + \mathfrak{s}^* \, \mathfrak{s}_a - 2 \mathfrak{s}_a \stackrel{[1]}{=} \Delta \, \mathfrak{s}_a \end{split}$$

Hence,  $[\mathfrak{s}_a, \Delta]\psi = 0$  and  $\mathfrak{s}_a$  commutes with  $\Delta$  [1].

(iv) We will see that the discrete Laplacian  $\Delta$  is selfadjoint: for all  $\varphi,\psi\in\ell^2(\mathbb{Z})$  we have

$$\begin{split} \left\langle \varphi, \Delta \psi \right\rangle &\stackrel{[1]}{=} \left\langle \varphi, \mathfrak{s}\psi \right\rangle + \left\langle \varphi, \mathfrak{s}^*\psi \right\rangle - 2 \left\langle \varphi, \psi \right\rangle \\ &\stackrel{[1]}{=} \left\langle \mathfrak{s}^*\varphi, \psi \right\rangle + \left\langle \mathfrak{s}\varphi, \psi \right\rangle - 2 \left\langle \varphi, \psi \right\rangle \\ &\stackrel{[1]}{=} \left\langle \Delta \varphi, \psi \right\rangle, \end{split}$$

i. e.  $\Delta^* = \Delta$  is selfadjoint [1].

(v) We apply  $\Delta$  to the sequence  $\psi_k$  with entries  $\psi_k(n) = e^{+ink}$ ,  $k \in [-\pi, +\pi]$  and obtain

$$(\Delta\psi_k)(n) = \psi_k(n+1) + \psi_k(n-1) - 2\psi_k(n) \stackrel{[1]}{=} \mathbf{e}^{+\mathbf{i}(n+1)k} + \mathbf{e}^{+\mathbf{i}(n-1)k} - 2\mathbf{e}^{+\mathbf{i}nk} \\ \stackrel{[1]}{=} (\mathbf{e}^{+\mathbf{i}k} + \mathbf{e}^{-\mathbf{i}k} - 2) \mathbf{e}^{+\mathbf{i}nk} \stackrel{[1]}{=} (2\cos k - 2) \mathbf{e}^{+\mathbf{i}nk} \stackrel{[1]}{=} E_k \psi_k(n).$$

Since  $|\psi_k(n)| = |\mathbf{e}^{+\mathbf{i}nk}| = 1$  is independent of  $n \in \mathbb{Z}$ , the sequence  $\psi_k$  cannot be square summable, because  $\psi_k \in \ell^2(\mathbb{Z})$  necessarily implies  $\lim_{|n|\to\infty} \psi_k(n) = 0$  [1].

# 14. Position and momentum representation (15 points)

Consider  $\ell^2(\mathbb{Z})$  with the usual scalar  $\langle \cdot, \cdot \rangle_{\ell^2(\mathbb{Z})}$  product and  $L^2([0, 2\pi])$  endowed with the scalar product

$$\left\langle \widehat{\varphi}, \widehat{\psi} \right\rangle_{L^2([0,2\pi])} := \frac{1}{2\pi} \int_0^{2\pi} \mathrm{d}k \, \overline{\widehat{\varphi}(k)} \, \widehat{\psi}(k).$$

Define the Fourier transform

$$\mathcal{F} : L^2([0, 2\pi]) \longrightarrow \ell^2(\mathbb{Z})$$
$$\psi(n) = \left(\mathcal{F}\widehat{\psi}\right)(n) := \left\langle \mathbf{e}^{+\mathbf{i}nk}, \widehat{\psi} \right\rangle_{L^2([0, 2\pi])}$$

and its inverse

$$\ell^2(\mathbb{Z}) \ni \psi \mapsto (\mathcal{F}^{-1}\psi)(k) = \sum_{n \in \mathbb{Z}} \psi(n) \, \mathbf{e}^{+\mathbf{i}nk}.$$

You may use without proof that  $\mathcal{F}$  is unitary.

- (i) For the shift operator  $(\mathfrak{s}\widehat{\psi})(n) := \widehat{\psi}(n-1)$ , compute  $\mathcal{F}^{-1}\mathfrak{s}\mathcal{F}$ .
- (ii) For the discrete Laplacian from problem 13, compute the momentum representation  $\mathcal{F}^{-1} \Delta \mathcal{F}$ .
- (iii) What is the connection between  $\psi_k$  from problem 13 (v) in the position representation and  $\mathcal{F}^{-1} \Delta \mathcal{F}$  in the momentum representation? Heuristically, what is the inverse Fourier transform of  $\psi_k$ ?
- (iv) Is  $\Delta \ge 0$ ? Justify your answer.

# Solution:

(i) Pick  $\widehat{\psi} \in L^2([0, 2\pi])$ . Then a straightforward computation yields

$$\begin{split} \left(\mathcal{F}^{-1}\,\mathfrak{s}\,\mathcal{F}\widehat{\psi}\right)(k) \stackrel{[1]}{=} \sum_{n\in\mathbb{Z}} \left(\mathfrak{s}\,\mathcal{F}\widehat{\psi}\right)(n)\,\mathbf{e}^{+\mathbf{i}nk} \stackrel{[1]}{=} \sum_{n\in\mathbb{Z}} (\mathcal{F}\widehat{\psi})(n-1)\,\mathbf{e}^{+\mathbf{i}k}\,\mathbf{e}^{+\mathbf{i}(n-1)k} \\ &= \mathbf{e}^{+\mathbf{i}k}\,\sum_{n'\in\mathbb{Z}} (\mathcal{F}\widehat{\psi})(n')\,\mathbf{e}^{+\mathbf{i}n'k} \stackrel{[1]}{=} \mathbf{e}^{+\mathbf{i}k}\,\widehat{\psi}(k). \end{split}$$

Hence, the discrete Laplacian in momentum representation is the multiplication operator  $\mathcal{F}^{-1}\mathfrak{s}\mathcal{F}=e^{+i\hat{k}}$  [1].

(ii) Given that  $\Delta = \mathfrak{s} + \mathfrak{s}^* - 2$  [1], we can reuse the result from (i) to conclude that in momentum representation, the discrete Laplacian is multiplication by  $2 \cos k - 2$ ,

$$\mathcal{F}^{-1} \Delta \mathcal{F} \stackrel{[1]}{=} \mathcal{F}^{-1} \left( \mathfrak{s} + \mathfrak{s}^* - 2 \right) \mathcal{F} \stackrel{[1]}{=} \mathbf{e}^{+\mathbf{i}\hat{k}} + \mathbf{e}^{-\mathbf{i}\hat{k}} - 2$$
$$\stackrel{[1]}{=} 2\cos\hat{k} - 2.$$

(iii) The  $\psi_k$  were pseudoeigenvectors to  $\Delta$  [1], meaning that while they satisfy the eigenvalue equation, they are not square-summable [1]. Heuristically, their inverse Fourier transform is the delta function  $\delta(\cdot - k)$  [1], because  $\mathcal{F}^{-1} \Delta \mathcal{F}$  is a multiplication operator and the eigenvectors of multiplication operators are delta distributions.

(iv) No,  $\Delta \leq 0$  : One can check that in Fourier representation, with  $\widehat{\psi} = \mathcal{F}^{-1}\psi$ 

$$\begin{split} \left\langle \psi, \Delta \psi \right\rangle_{\ell^{2}(\mathbb{Z})} &\stackrel{[1]}{=} \left\langle \mathcal{F}^{-1}\psi, \mathcal{F}^{-1}\Delta \mathcal{F}\mathcal{F}^{-1}\psi \right\rangle_{L^{2}([0,2\pi])} \\ &\stackrel{[1]}{=} \left\langle \widehat{\psi}, \left(2\cos\hat{k}-2\right)\widehat{\psi} \right\rangle_{L^{2}([0,2\pi])} \\ &\stackrel{[1]}{=} \frac{1}{2\pi} \int_{0}^{2\pi} \mathrm{d}k \underbrace{\left(2\cos k-2\right)}_{\leq 0} \left|\widehat{\psi}(k)\right|^{2} \stackrel{[1]}{\leq} 0 \end{split}$$

#### 15. Rank-1 operators (14 points)

Suppose  $\varphi, \psi \neq 0$  are elements of a Hilbert space  $\mathcal{H}$ , and define the rank-1 operator  $T = |\varphi\rangle\langle\psi|$  via

$$T\phi = \langle \psi, \phi \rangle \varphi.$$

- (i) Find all eigenvectors and eigenvalues of T.
- (ii) Compute  $\sigma(T)$ .
- (iii) Determine the nature of the spectrum, i. e. determine  $\sigma_p(T)$ ,  $\sigma_{cont}(T)$  and  $\sigma_r(T)$ .

#### Solution:

(i) We can read off the eigenvalues from the form of the operator: the first eigenvector is  $\varphi$  [1] with eigenvalue  $\lambda := \langle \psi, \varphi \rangle$  [1].

The other eigenvalue is 0 [1], because for any vector  $\phi$  perpendicular to  $\psi$ , we have  $T\phi = 0$  [1], and thus the eigenspace is

$$\ker T \stackrel{[1]}{=} \{\psi\}^{\perp}.$$

Now there are two cases:  $\psi \perp \varphi$ , and then also  $\lambda = 0$  and the only eigenvalue is 0 [1]. Or  $\langle \psi, \varphi \rangle \neq 0$  and T has two different eigenvalues [1].

(ii) Clearly,  $\{0, \lambda\} \subseteq \sigma(T)$  where  $\lambda = \langle \psi, \varphi \rangle$  [1].

Since ran  $T = \text{span}\{\varphi\}$  is a one-dimensional subspace, the operator T - z is always invertible on  $(\operatorname{ran} T)^{\perp}$  [1]. On the one-dimensional subspace ran T, the operator is invertible if and only if  $z \neq \lambda$  [1]. Hence, we have shown  $\sigma(T) = \{0, \lambda\}$  [1].

(iii) By the classification introduced in Chapter 4.1, we know that

$$\sigma_{\mathbf{p}}(T) = \sigma(P) \stackrel{[1]}{=} \{0, \lambda\}, \qquad \qquad \sigma_{\mathsf{cont}}(T) \stackrel{[1]}{=} \emptyset, \qquad \qquad \sigma_{\mathbf{r}}(T) \stackrel{[1]}{=} \emptyset.$$