



Unitary Evolution Group, Resolvents & Symmetric Operators

Homework Problems

16. Translation semigroup on $L^2([0, +\infty))$ (24 points)

(i) Show that for $t \geq 0$, the translation operator

$$(T_t \psi)(x) := \begin{cases} 0 & x \in [0, t) \\ \psi(x - t) & x \in [t, +\infty) \end{cases}$$

preserves angles on $L^2([0, +\infty))$, i. e. $\langle T_t \psi, T_t \varphi \rangle = \langle \psi, \varphi \rangle$.

(ii) Compute the adjoint of T_t .

(iii) Show that $\{T_t\}_{t \in [0, +\infty)}$ forms a *semigroup*, i. e. $T_{t_1} T_{t_2} = T_{t_1+t_2}$ holds for all $t_1, t_2 \in [0, +\infty)$ and $T_0 = \text{id}$.

(iv) Find the generator of $\{T_t\}_{t \in [0, +\infty)}$.

(A formal computation ignoring domain questions suffices.)

(v) Is the generator of $\{T_t\}_{t \in [0, +\infty)}$ symmetric on $C_c^\infty([0, +\infty))$? Justify your answer.

(vi) Find a domain such that the generator of $\{T_t\}_{t \in [0, +\infty)}$ is symmetric.

(vii) Can $\{T_t\}_{t \in [0, +\infty)}$ be extended to a unitary evolution group? Justify your answer.

Solution:

(i) For $\psi, \varphi \in L^2([0, +\infty))$ we compute

$$\begin{aligned} \langle T_t \psi, T_t \varphi \rangle &\stackrel{[1]}{=} \int_0^{+\infty} dx \overline{(T_t \psi)(x)} (T_t \varphi)(x) \stackrel{[1]}{=} \int_0^t dx 0 + \int_t^{+\infty} dx \overline{\psi(x-t)} \varphi(x-t) \\ &\stackrel{[1]}{=} \int_0^{+\infty} dx \overline{\psi(x)} \varphi(x) \stackrel{[1]}{=} \langle \psi, \varphi \rangle. \end{aligned}$$

(ii) A quick computation reveals that the adjoint operator is $(T_y^* \psi)(x) = \psi(x + y)$ [1]:

$$\begin{aligned} \langle T_t \psi, \varphi \rangle &\stackrel{[1]}{=} \int_0^{+\infty} dx \overline{(T_t \psi)(x)} \varphi(x) \\ &\stackrel{[1]}{=} \int_0^t dx 0 + \int_t^{+\infty} dx \overline{\psi(x-t)} \varphi(x) \\ &\stackrel{[1]}{=} \int_0^{+\infty} dx \overline{\psi(x)} \varphi(x+t) \stackrel{[1]}{=} \langle \psi, T_t^* \varphi \rangle \end{aligned}$$

(iii) That $T_0 = \text{id}$ is clear from the definition. Pick $t_1, t_2 \in [0, +\infty)$ and $\psi \in L^2([0, +\infty))$. Then we get from the definition

$$\begin{aligned} (T_{t_1} T_{t_2} \psi)(x) &\stackrel{[1]}{=} \begin{cases} 0 & x \in [0, t_1) \\ (T_{t_2} \psi)(x - t_1) & x \in [t_1, +\infty) \end{cases} \\ &\stackrel{[1]}{=} \begin{cases} 0 & x \in [0, t_1) \\ 0 & x \in [t_1, t_2) \\ \psi(x - t_1 - t_2) & x \in [t_2, +\infty) \end{cases} \\ &\stackrel{[1]}{=} (T_{t_1+t_2} \psi)(x). \end{aligned}$$

(iv) For the purpose of the computation we may assume $x > 0$. The formally, we obtain

$$\begin{aligned} i \frac{d}{dt} (T_t \psi)(x) \Big|_{t=0} &\stackrel{[1]}{=} i \partial_t (\psi(x - t)) \Big|_{t=0} \stackrel{[1]}{=} -i \partial_x \psi(x - t) \Big|_{t=0} \\ &\stackrel{[1]}{=} -i \partial_x \psi(x). \end{aligned}$$

(v) The generator $-i \partial_x$ is not symmetric on $\mathcal{C}_c^\infty([0, +\infty))$, because the boundary terms do not vanish [1]:

$$\begin{aligned} \langle -i \partial_x \psi, \varphi \rangle &\stackrel{[1]}{=} +i \int_0^{+\infty} dx \overline{\partial_x \psi(x)} \varphi(x) \\ &\stackrel{[1]}{=} +i \left[\overline{\psi(x)} \varphi(x) \right]_0^{+\infty} - i \int_0^{+\infty} dx \overline{\psi(x)} \partial_x \varphi(x) \\ &\stackrel{[1]}{=} -i \overline{\psi(0)} \varphi(0) + \langle \psi, -i \partial_x \varphi \rangle \end{aligned}$$

(vi) We need to ensure $\psi(0) = 0$ [1], so $\mathcal{D} := \{ \psi \in \mathcal{C}_c^\infty([0, +\infty)) \mid \psi(0) = 0 \}$ will do [1].

(vii) No, the adjoint $(T_t^* \psi) = \psi(x + t)$ is not norm-preserving [1], and hence, not unitary [1]. (For instance, if we fix $t > 0$ and pick $0 \neq \varphi \in L^2([0, +\infty))$ so that $\varphi(x)$ is 0 almost everywhere as long as $x \geq t$, then $T_t^* \varphi = 0$.) Thus, $\{T_t\}_{t \in [0, +\infty)}$ cannot be extended to a unitary evolution group [1].

17. Convergence of operators (20 points)

Consider the following sequences $\{T_n\}_{n \in \mathbb{N}}$ of operators on the Hilbert space

$$\ell^2(\mathbb{N}) = \left\{ a \equiv (a_n)_{n \in \mathbb{N}} \mid \sum_{n=1}^{\infty} |a_n|^2 < \infty \right\}$$

and investigate whether they converge in norm, strongly or weakly:

- (i) $T_n(a) := (\frac{1}{n}a_1, \frac{1}{n}a_2, \dots)$
- (ii) $T_n(a) := (\underbrace{0, \dots, 0}_{n \text{ places}}, a_{n+1}, a_{n+2}, \dots)$
- (iii) $T_n(a) := (\underbrace{0, \dots, 0}_{n \text{ places}}, a_1, a_2, \dots)$

Solution:

- (i) The sequence T_n converges in norm/uniformly to $0 \in \mathcal{B}(\ell^2(\mathbb{N}))$ [1], because

$$\|T_n(a)\|_{\ell^2(\mathbb{N})} = \frac{1}{n} \|a\|_{\ell^2(\mathbb{N})} \xrightarrow{n \rightarrow \infty} 0, \quad [1]$$

and thus $\|T_n\|_{\mathcal{B}(\ell^2(\mathbb{N}))} = 1/n$. The above equation also implies that T_n converges to 0 also strongly and weakly [2], because

$$|\langle a, T_n(b) \rangle_{\ell^2(\mathbb{N})}| \leq \|a\|_{\ell^2(\mathbb{N})} \|T_n(b)\|_{\ell^2(\mathbb{N})} \xrightarrow{n \rightarrow \infty} 0. \quad [1]$$

- (ii) For a fixed $a \in \ell^2(\mathbb{N})$, we have

$$\|T_n(a)\|_{\ell^2(\mathbb{N})} \stackrel{[1]}{=} \sum_{j=1}^{\infty} |(T_n(a))_j|^2 = \sum_{j=n+1}^{\infty} |a_j|^2 \xrightarrow{n \rightarrow \infty} 0, \quad [1]$$

and thus T_n converges strongly (and weakly) to $0 \in \mathcal{B}(\ell^2(\mathbb{N}))$ [2]. However, if $e_n := (\delta_{jn})_{j \in \mathbb{N}} = (0, \dots, 0, 1, 0, \dots)$, we see that

$$\|T_n e_{n+1}\|_{\ell^2(\mathbb{N})} \stackrel{[1]}{=} 1,$$

and thus T_n does not converge to 0 in norm [1], because $\|T_n\|_{\mathcal{B}(\ell^2(\mathbb{N}))} \geq 1$ [1].

- (iii) T_n converges weakly to 0 [1]:

$$\begin{aligned} |\langle a, T_n(b) \rangle_{\ell^2(\mathbb{N})}| &\stackrel{[1]}{=} \left| \sum_{j=1}^{\infty} \bar{b}_j (T_n(a))_j \right| \stackrel{[1]}{=} \left| \sum_{j=n+1}^{\infty} \bar{b}_j a_{j-n} \right| \\ &\leq \left(\sum_{j=n+1}^{\infty} |b_j|^2 \right)^{1/2} \left(\sum_{j=n+1}^{\infty} |a_{j-n}|^2 \right)^{1/2} \stackrel{[1]}{=} \left(\sum_{j=n+1}^{\infty} |b_j|^2 \right)^{1/2} \left(\sum_{j=1}^{\infty} |a_j|^2 \right)^{1/2} \\ &\xrightarrow{n \rightarrow \infty} 0. \quad [1] \end{aligned}$$

However, it does not converge strongly or in norm, because

$$\|T_n(a)\|_{\ell^2(\mathbb{N})}^2 \stackrel{[1]}{=} \sum_{j=n+1}^{\infty} |a_{j-n}|^2 \stackrel{[1]}{=} \sum_{j=1}^{\infty} |a_j|^2 \stackrel{[1]}{=} \|a\|_{\ell^2(\mathbb{N})}^2.$$

18. The resolvent (27 points)

Let $S, T \in \mathcal{B}(\mathcal{X})$ be operators on a Banach space \mathcal{X} with resolvent sets $\rho(S)$ and $\rho(T)$. On these sets, the resolvents $(T - z)^{-1}$ and $(S - z)^{-1}$ exist as bounded operators.

(i) Prove the first resolvent identity, i. e. that for any $z, z' \in \rho(T)$ we have

$$(T - z)^{-1} - (T - z')^{-1} = (z - z') (T - z)^{-1} (T - z')^{-1}.$$

(ii) Prove the second resolvent identity, i. e. that for any $z \in \rho(T) \cap \rho(S)$ we have

$$(T - z)^{-1} - (S - z)^{-1} = (T - z)^{-1} (T - S) (S - z)^{-1}.$$

(iii) Prove that if $\|T\| < 1$, then the geometric series $\sum_{n=0}^{\infty} T^n$ exists in $\mathcal{B}(\mathcal{X})$ and equals $(\text{id} - T)^{-1}$.

(iv) Show that the resolvent set $\rho(T) \subseteq \mathbb{C}$ is open and the resolvent is $z \mapsto (T - z)^{-1}$ is analytic on $\rho(T)$, meaning locally there exists a power series expansion of $(T - z)^{-1}$ which converges in operator norm.

(v) Show that the spectrum $\sigma(T) \subseteq \mathbb{C}$ is closed.

Solution:

(i) If we multiply the left-hand side with $T - z$ from the left and $T - z'$ from the right [1], we get

$$(T - z) ((T - z)^{-1} - (T - z')^{-1}) (T - z') \stackrel{[1]}{=} (T - z') - (T - z) \stackrel{[1]}{=} z - z'.$$

Evidently, this is equivalent to the first resolvent identity.

(ii) We again multiply with $T - z$ from the left and $S - z$ from the right [1], and obtain

$$(T - z) ((T - z)^{-1} - (S - z)^{-1}) (S - z) \stackrel{[1]}{=} (S - z) - (T - z) \stackrel{[1]}{=} S - T.$$

(iii) First of all, the sequence of partial sums $S_N := \sum_{n=0}^N T^n$ is Cauchy [1], because for $N > M$

$$\|S_N - S_M\| = \left\| \sum_{n=M+1}^N T^n \right\| \stackrel{[1]}{\leq} \sum_{n=M+1}^N \|T\|^n$$

and the right-hand side is finite since $\|T\| < 1$ [1], and goes to 0 as $N, M \rightarrow \infty$ [1]. By the completeness of $\mathcal{B}(\mathcal{X})$ (Proposition 4.1.4) [1], the sequence of partial sums S_N converges to some $S \in \mathcal{B}(\mathcal{X})$ [1]. To show $S = (\text{id} - T)^{-1}$, we compute

$$\begin{aligned} (\text{id} - T) S &\stackrel{[1]}{=} \lim_{N \rightarrow \infty} (\text{id} - T) \sum_{n=0}^N T^n \\ &\stackrel{[1]}{=} \lim_{N \rightarrow \infty} (\text{id} - T^{N+1}) \stackrel{[1]}{=} \text{id}. \end{aligned}$$

(iv) Let $z_0 \in \rho(T)$. We first show that there exists an $\varepsilon > 0$ such that $\{z \in \mathbb{C} \mid |z - z_0| < \varepsilon\} \subseteq \rho(T)$.

Pick $0 < \varepsilon < \|(T - z_0)^{-1}\|^{-1}$ [1]. Then

$$(T - z) \stackrel{[1]}{=} (T - z_0) \left(\text{id} - (z - z_0) (T - z_0)^{-1} \right)$$

is invertible [1], because $(T - z_0)^{-1}$ exists and in view of

$$\left\| (z - z_0) (T - z_0)^{-1} \right\| < \varepsilon \left\| (T - z_0)^{-1} \right\| \stackrel{[1]}{<} 1$$

(iii) tells us that also the second term in the product is invertible [1]. That means $\rho(T)$ is open [1].

(iii) also gives us an explicit expression for the power series of the resolvent in this neighborhood,

$$(T - z)^{-1} \stackrel{[1]}{=} \left(\text{id} - (z - z_0) (T - z_0)^{-1} \right)^{-1} (T - z_0)^{-1} \stackrel{[1]}{=} \sum_{n=0}^{\infty} (z - z_0)^n (T - z_0)^{-(n+1)}.$$

Hence, $(T - z)^{-1}$ is analytic in a vicinity of z_0 [1], and since the point was chosen arbitrarily, the resolvent is in fact analytic on all of $\rho(T)$ [1].

(v) Since $\rho(T)$ is open by (iv) [1], its complement in \mathbb{C} , the spectrum $\sigma(T)$, is closed [1].

19. Symmetric operators (17 points)

Let $H = \frac{1}{2m}(-i\nabla_x)^2 + V$ be a Hamilton operator with potential $V \in \mathcal{C}(\mathbb{R}^3, \mathbb{R})$.

Define the smooth functions with compact support as

$$\mathcal{C}_c^\infty(\mathbb{R}^3) := \{\varphi : \mathbb{R}^3 \longrightarrow \mathbb{C} \mid \varphi \in \mathcal{C}^\infty(\mathbb{R}^3), \text{ supp } \varphi \text{ compact}\}.$$

- (i) Prove $\mathcal{C}_c^\infty(\mathbb{R}^3) \subset L^2(\mathbb{R}^3)$.
(ii) Show that H is symmetric on $\mathcal{C}_c^\infty(\mathbb{R}^3)$, i. e. that

$$\langle \varphi, H\psi \rangle = \langle H\varphi, \psi \rangle$$

holds for all $\varphi, \psi \in \mathcal{C}_c^\infty(\mathbb{R}^3)$.

Solution:

- (i) Every smooth function with compact support is square-integrable: let $\varphi \in \mathcal{C}_c^\infty(\mathbb{R}^3)$, then there exists a compact subset $K \subset \mathbb{R}^3$, so that

$$\text{supp } \varphi = \overline{\{x \in \mathbb{R}^3 \mid \varphi(x) \neq 0\}} \subseteq K. \quad [1]$$

Since φ is also continuous, we can estimate the supremum from above by

$$\|\varphi\|_\infty = \sup_{x \in \mathbb{R}^3} |\varphi(x)| = \sup_{x \in K} |\varphi(x)| < \infty. \quad [1]$$

Hence, we obtain

$$\|\varphi\|^2 = \int_{\mathbb{R}^3} dx |\varphi(x)|^2 \stackrel{[1]}{=} \int_K dx |\varphi(x)|^2 \stackrel{[1]}{\leq} |K| \left(\sup_{x \in K} |\varphi(x)|\right)^2 < \infty.$$

- (ii) We will treat kinetic and potential energy separately: clearly, derivatives map $\mathcal{C}_c^\infty(\mathbb{R}^3)$ into itself, and thus $(-i\nabla_x)^2\varphi \in L^2(\mathbb{R}^3)$ [1]. Fix $\varphi, \psi \in \mathcal{C}_c^\infty(\mathbb{R}^3)$. Then there exists a compact set $K \subset \mathbb{R}^3$ whose interior contains $\text{supp } \varphi$ and $\text{supp } \psi$ [1]. Then we compute using repeated partial integration

$$\begin{aligned} \langle \varphi, \frac{1}{2m}(-i\nabla_x)^2\psi \rangle &= \sum_{j=1}^3 \frac{1}{2m} \langle \varphi, (-i\partial_{x_j})^2\psi \rangle \stackrel{[1]}{=} \sum_{j=1}^3 \frac{1}{2m} \int_{\mathbb{R}^3} dx \overline{\varphi(x)} ((-i\partial_{x_j})^2\psi)(x) \\ &\stackrel{[1]}{=} \sum_{j=1}^3 \frac{1}{2m} \int_K dx \overline{\varphi(x)} ((-i\partial_{x_j})^2\psi)(x) \\ &\stackrel{[1]}{=} \sum_{j=1}^3 \frac{1}{2m} \int_{\partial K} dS(x) \overline{\varphi(x)} ((-i)^2\partial_{x_j}\psi)(x) + \\ &\quad - \sum_{j=1}^3 \frac{1}{2m} \int_K dx \overline{\partial_{x_j}\varphi(x)} ((-i)^2\partial_{x_j}\psi)(x) \\ &\stackrel{[1]}{=} 0 - \sum_{j=1}^3 \frac{1}{2m} (-i)^2 \int_{\partial K} dS(x) \overline{\partial_{x_j}\varphi(x)} \psi(x) + \\ &\quad + \sum_{j=1}^3 \frac{1}{2m} (-i)^2 \int_K dx \overline{\partial_{x_j}^2\varphi(x)} \psi(x) \end{aligned}$$

$$\begin{aligned} &\stackrel{[1]}{=} \sum_{j=1}^3 \frac{1}{2m} \int_K \mathbf{d}x \overline{((-i\partial_{x_j})^2 \varphi)(x)} \psi(x) \\ &\stackrel{[1]}{=} \left\langle \frac{1}{2m} (-i\nabla_x)^2 \varphi, \psi \right\rangle \end{aligned}$$

Here, $dS(x)$ is the surface measure on ∂K . The boundary terms vanish, because φ and ψ as well as their derivatives vanish on ∂K .

Now to the potential energy: since V is continuous, it is bounded on compact subsets. Choose any $\varphi, \psi \in C_c^\infty(\mathbb{R}^3)$. Then $V\varphi \in L^2(\mathbb{R}^3)$ [1] and hence,

$$\begin{aligned} \langle \varphi, V\psi \rangle &\stackrel{[1]}{=} \int_{\mathbb{R}^3} \mathbf{d}x \overline{\varphi(x)} (V\psi)(x) \stackrel{[1]}{=} \int_{\mathbb{R}^3} \mathbf{d}x \overline{\varphi(x)} V(x) \psi(x) \\ &\stackrel{[1]}{=} \int_{\mathbb{R}^3} \mathbf{d}x \overline{(V\varphi)(x)} \psi(x) \stackrel{[1]}{=} \langle V\varphi, \psi \rangle \end{aligned}$$

holds.